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# Wrinkling of a stiff thin film bonded to a pre-strained, compliant substrate with finite thickness

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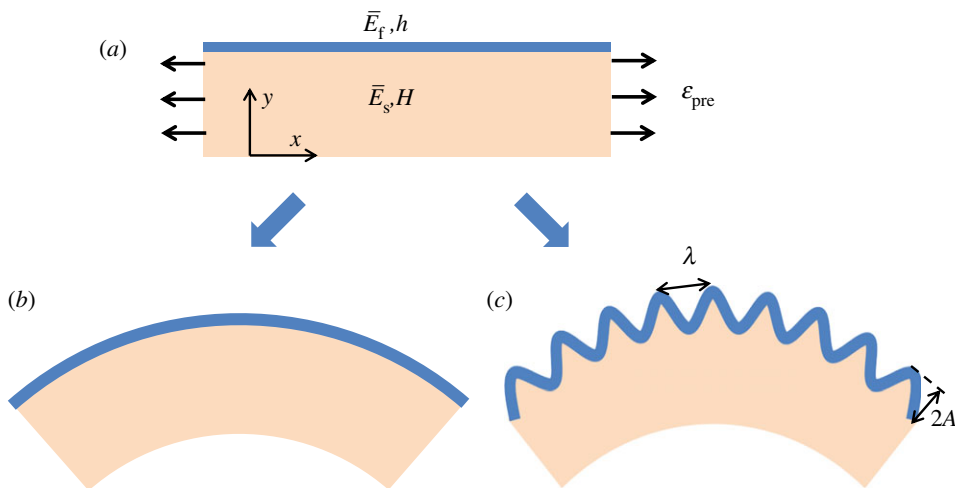
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A stiff thin film bonded to a pre-strained, compliant substrate wrinkles into a sinusoidal form upon release of the pre-strain. Many analytical models developed for the critical pre-strain for wrinkling assume that the substrate is semi-infinite. This critical pre-strain is actually much smaller than that for a substrate with finite thickness (Ma Y *et al.* 2016 *Adv. Funct. Mater.* (doi:10.1002/adfm.201600713)). An analytical solution of the critical pre-strain for a system of a stiff film bonded to a pre-strained, finite-thickness, compliant substrate is obtained, and it agrees well with the finite-element analysis. The finite-thickness effect is significant when the substrate tensile stiffness cannot overwhelm the film tensile stiffness.

## 1. Introduction

A stiff film bonded to a pre-strained, compliant substrate wrinkles upon releasing the pre-strain [1,2]. Such a system has many important applications in stretchable



**Figure 1.** Schematic illustrations. (a) A stiff thin film bonded to a pre-strained, compliant substrate with finite thickness; (b) bending of the film/substrate system upon release of the small pre-strain; and (c) wrinkling of the stiff thin film, along with bending of the film/substrate system, upon release of large pre-strain. (Online version in colour.)

inorganic electronics [3–8], micro/nano pattern formation [9–11], high-precision micro/nano measurement techniques [12], tuneable metamaterials [13], nanocomposites [14], stretchable transistors [15] and biomimetic materials [16]. Analytical models have been developed for wrinkling of a stiff thin film on a pre-strained compliant substrate [17–21]. The results identify the critical pre-strain for wrinkling, below which the film remains flat. However, all of these studies assume that the substrate is semi-infinite such that its tensile stiffness overwhelms that of the film. Consequently, the substrate recovers the initial length after the pre-strain is released and its bottom remains flat.<sup>1</sup>

The critical pre-strain for wrinkling obtained for a semi-infinite substrate, however, is smaller than the numerical and experimental results for a substrate with finite thickness [22], even for substrates that are more than 1000 times thicker than the film. This is because the substrate elastic modulus  $E_s$  is often more than five orders of magnitude smaller than the film elastic modulus  $E_f$  [1,2], such that the substrate tensile stiffness  $E_s H$  cannot overwhelm the film tensile stiffness  $E_f h$ , where  $H$  and  $h$  are the substrate and film thicknesses, respectively (figure 1a). Consequently,

- (1) the substrate cannot shrink back to its initial length after release of the pre-strain; and
- (2) the film/substrate system may bend after the pre-strain is released (figure 1b).

The recent study by Ma *et al.* [22] accounted for (1), while this paper aims to establish an analytic model for both (1) and (2). The resulting critical pre-strain will be useful for many applications such as the strain-limiting design of materials [22] and tuneable optical design of the intensity for diffraction peaks [23].

## 2. Analytical model

A stiff thin film is bonded onto a pre-strained ( $\epsilon_{\text{pre}}$ ), compliant substrate (figure 1a). For small pre-strain, the stiff film does not wrinkle upon release of the pre-strain; instead, the film and substrate bend (figure 1b). Let  $\epsilon$  denote the membrane strain in the film. The strain in the substrate

<sup>1</sup>Huang *et al.* [24] studied the finite thickness of a substrate subjected to compression, to be discussed below.

is  $\varepsilon_s(y) = \varepsilon_{\text{pre}} + \varepsilon - \kappa(H - y)$ , where  $\kappa$  is the curvature of the substrate, and the co-ordinate  $y$  is shown in figure 1a. The potential energy is

$$U_{\text{bend}} = \frac{1}{2} \bar{E}_f h \varepsilon^2 + \frac{1}{2} \bar{E}_s \int_0^H [\varepsilon_{\text{pre}} + \varepsilon - \kappa(H - y)]^2 dy, \quad (2.1)$$

where  $\bar{E}_f = E_f/(1 - \nu_f^2)$  and  $\bar{E}_s = E_s/(1 - \nu_s^2)$  are the plane-strain moduli of the stiff thin film and compliant substrate, respectively, and  $\nu_f$  and  $\nu_s$  are the Poisson's ratios.

Minimization of the potential energy  $\partial U_{\text{bend}}/\partial \varepsilon = 0$  and  $\partial U_{\text{bend}}/\partial \kappa = 0$  gives  $\varepsilon = -\bar{E}_s H \varepsilon_{\text{pre}} / (4\bar{E}_f h + \bar{E}_s H)$  and  $\kappa = -6\bar{E}_f h \varepsilon_{\text{pre}} / [H(4\bar{E}_f h + \bar{E}_s H)]$ . Equation (2.1) then becomes

$$U_{\text{bend}} = \frac{\bar{E}_f h \bar{E}_s H \varepsilon_{\text{pre}}^2}{2(4\bar{E}_f h + \bar{E}_s H)}. \quad (2.2)$$

Once the pre-strain exceeds the critical pre-strain (to be determined), the stiff film wrinkles on the top surface of the substrate (figure 1c) and the film/substrate bends. In addition to the membrane strain  $\varepsilon$ , the film is also subjected to wrinkling with amplitude  $A$  and period  $\lambda$  to be determined. The strain energy in the film is [24]

$$U_{\text{film}} = \frac{1}{2} \bar{E}_f h \varepsilon^2 - \frac{1}{4} \bar{E}_f h |\varepsilon| k^2 A^2 + \frac{1}{32} \bar{E}_f h k^4 A^4 + \frac{1}{48} \bar{E}_f h^3 k^4 A^2, \quad (2.3)$$

which degenerates to the first term on the right-hand side of equation (2.1) when the amplitude  $A = 0$ ; here  $k = 2\pi/\lambda$ . The strain energy in the substrate is

$$U_{\text{substrate}} = \frac{1}{2} \bar{E}_s \int_0^H [\varepsilon_{\text{pre}} + \varepsilon - \kappa(H - y)]^2 dy + \frac{\bar{E}_s}{4} k A^2 g(kH), \quad (2.4)$$

which degenerates to the last term in equation (2.1) when the amplitude  $A = 0$ . The last term in the above equation is the strain energy in the substrate due to wrinkling [24], and the function  $g$  is

$$g(x) = \frac{\cosh(2x) + 1 + 2x^2}{2 \sinh(2x) - 4x}, \quad (2.5)$$

for an incompressible substrate ( $\nu_s = 0.5$ ), and  $g$  approaches  $1/2$  for a semi-infinite substrate. The potential energy is the sum of  $U_{\text{film}}$  and  $U_{\text{substrate}}$ ,

$$U_{\text{bend+wrinkle}} = \frac{1}{2} \bar{E}_f h \varepsilon^2 + \frac{1}{2} \bar{E}_s \int_0^H [\varepsilon_{\text{pre}} + \varepsilon - \kappa(H - y)]^2 dy + \frac{1}{4} \bar{E}_f h (f - |\varepsilon|) k^2 A^2 + \frac{1}{32} \bar{E}_f h k^4 A^4, \quad (2.6)$$

where

$$f = \frac{k^2 h^2}{12} + \frac{\bar{E}_s g(kH)}{kh \bar{E}_f}. \quad (2.7)$$

Minimization of the potential energy with respect to  $k$  and  $A$ ,  $\partial U_{\text{bend+wrinkle}}/\partial k = 0$  and  $\partial U_{\text{bend+wrinkle}}/\partial A = 0$ , gives

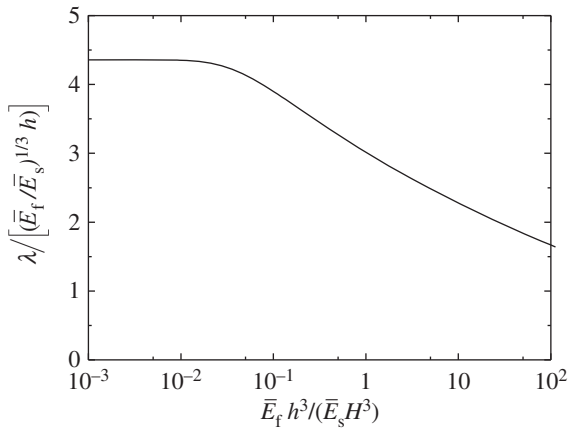
$$6 \frac{g(kH) - g'(kH)kH}{(kH)^3} = \frac{\bar{E}_f h^3}{\bar{E}_s H^3} \quad (2.8)$$

and

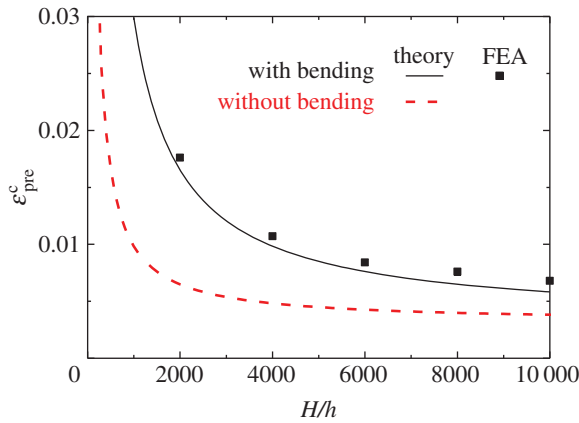
$$k^2 A^2 = 4(|\varepsilon| - f), \quad (2.9)$$

where  $g'(x) = dg(x)/dx$ . Equation (2.8) suggests that the normalized period,  $\lambda/[(\bar{E}_f/\bar{E}_s)^{1/3}h]$ , or equivalently  $kh/(\bar{E}_s/\bar{E}_f)^{1/3}$ , depends only on the film/substrate bending stiffness ratio  $\bar{E}_f h^3/\bar{E}_s H^3$ , as shown in figure 2. The period becomes independent of the substrate thickness  $H$  when the bending stiffness ratio  $\bar{E}_f h^3/\bar{E}_s H^3$  is less than 0.01, which is consistent with Huang *et al.* [24].

Minimization of the potential energy with respect to  $\varepsilon$  and  $\kappa$ ,  $\partial U_{\text{bend+wrinkle}}/\partial \varepsilon = 0$  and  $\partial U_{\text{bend+wrinkle}}/\partial \kappa = 0$ , gives  $\varepsilon = 4\bar{E}_f h f / (\bar{E}_s H) - \varepsilon_{\text{pre}}$ , and  $\kappa = 6\bar{E}_f h f / (\bar{E}_s H^2)$ , where  $f$  is obtained



**Figure 2.** The normalized wrinkle period  $\lambda / [(\bar{E}_f / \bar{E}_s)^{1/3} h]$  versus the film-to-substrate bending stiffness ratio  $[\bar{E}_f h^3 / (\bar{E}_s H^3)]$ .



**Figure 3.** The critical pre-strain  $\varepsilon_{pre}^c$  versus the substrate-to-film thickness ratio ( $H/h$ ) for a polyimide film on a PDMS substrate. FEA, finite-element analysis; PDMS, polydimethylsiloxane. (Online version in colour.)

from equation (2.7). The potential energy then becomes

$$U_{\text{bend+wrinkle}} = \bar{E}_f h f \left[ \varepsilon_{\text{pre}} - \frac{1}{2} \left( \frac{4\bar{E}_f h}{\bar{E}_s H} + 1 \right) f \right]. \quad (2.10)$$

Comparison of the potential energy in equations (2.2) and (2.10) suggests that wrinkling occurs when  $U_{\text{bend}} > U_{\text{bend+wrinkle}}$ , which gives

$$\varepsilon_{\text{pre}} > \left( \frac{4\bar{E}_f h}{\bar{E}_s H} + 1 \right) f = \left( \frac{4\bar{E}_f h}{\bar{E}_s H} + 1 \right) \left[ \frac{k^2 h^2}{12} + \frac{\bar{E}_s g(kH)}{kh\bar{E}_f} \right], \quad (2.11)$$

where  $k, f$  and  $g$  are obtained from equations (2.5), (2.7) and (2.8), respectively. It should be pointed out that equation (2.11) also ensures that the right-hand side of equation (2.9) is positive such that there is a solution for the amplitude  $A$ .

### 3. Discussion

When the bending stiffness of the substrate overwhelms that of the film, i.e.  $\bar{E}_f h^3 / (\bar{E}_s H^3) < \sim 0.01$ , equation (2.11) can be further simplified as

$$\varepsilon_{\text{pre}} > \varepsilon_{\text{pre}}^c = \frac{1}{4} \left( \frac{4\bar{E}_f h}{\bar{E}_s H} + 1 \right) \left( \frac{3\bar{E}_s}{\bar{E}_f} \right)^{2/3}. \quad (3.1)$$

For  $H \rightarrow \infty$ , the above equation degenerates to that for a semi-infinite substrate [1,2]. The critical pre-strain  $\varepsilon_{\text{pre}}^c$  in equation (3.1) is larger than  $\frac{1}{4}((\bar{E}_f h / \bar{E}_s H) + 1)(3\bar{E}_s / \bar{E}_f)^{2/3}$  [22], which neglects the effect of film/substrate bending. Figure 3 shows the critical pre-strain  $\varepsilon_{\text{pre}}^c$  versus the thickness ratio  $H/h$  for a polyimide film ( $E_f = 2.5$  GPa,  $\nu_f = 0.34$ ) on a polydimethylsiloxane substrate ( $E_s = 1$  MPa,  $\nu_s = 0.5$ ). The results obtained from finite-element analysis agree well with the critical pre-strain in equation (3.1).

**Competing interests.** We declare we have no competing interests.

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